

# CO Dissociation and CO + O Reactions on a Nanosized Iron Cluster

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## ABSTRACT

Catalysis over metal nanoparticles is essential for the growth of carbon nanotubes and all the properties of the resulting nanotube, such as diameter and chirality, are affected by the metal particle. Thus, it is very important to understand the carbon chemistry taking place on nanometer size metal particles. Here we present the first ab initio computational study of chemical reactions on a nanosized iron cluster. The clusters have reaction sites, such as edges and vertexes between the facets, Which have not been studied before. First principles electronic structure calculations, fully incorporating the effects of spin polarization and non-collinear magnetic moments, have been used to investigate CO disproportionation on an isolated Fe<sub>55</sub> cluster. After CO dissociation, O atoms remain on the surface while C atoms move into the cluster, presumably as the initial step toward carbide formation. Here we show that the lowest CO dissociation barrier found on the cluster (0.77 eV) is lower than on most previously studied Fe surfaces. This dissociation occurs on a vertex between the facets. Several possible paths for CO<sub>2</sub> formation were identified. The calculated lowest reaction barrier is 1.08 eV, which is comparable to the barrier of 0.65 eV obtained by experiment.

## KEYWORDS

Nanoparticle, density functional theory (DFT), mechanism, synthesis

Since their discovery, carbon nanotubes (CNTs) have been the object of very intensive research. CNTs are of great interest since they exhibit unique and useful chemical and physical properties related to toughness, electrical/thermal conductivity, and magnetism [1]. Several different techniques have been introduced to synthesize CNTs [2–5] but among the various routes, carbon vapor deposition (CVD) methods have attracted widespread interest since they open a way to highly controlled and continuous CNT production on designed architectures [5–9]. In

these processes, metal nanoparticles are produced in a mixed flow of carbon precursors and other gases (e.g., hydrogen) and the growth process is driven by cleaving the carbon atoms from the precursors and these atoms then form carbon structures on the surface of the nanoparticles. The metal catalyst nanoparticle is essential for CNT growth. All properties, like diameter and chirality, of the nanotube are determined by the metal particle. The ability to control the properties of CNTs is essential for their use in technological applications [10]. Thus, it

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is very important to understand the carbon chemistry taking place on the nanometer size metal particles.

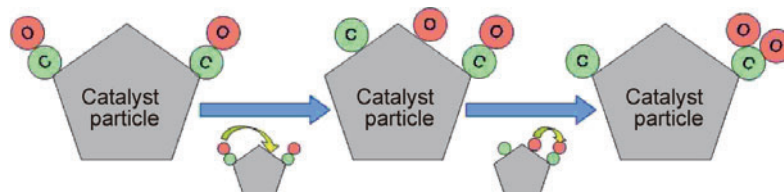
In addition to the CNT synthesis, metal nanoclusters with a size less than 10 nm have attracted a great deal of attention due to their applications in magnetism [11], electronics [12], and catalysis [13]. Often the good catalytic activity can be related to catalytic sites, like atomic size steps, on the cluster. Due to the high curvature of the clusters, the special site density and distribution are much higher than on almost all other surfaces. Furthermore, real nanoclusters have several unique active sites like facets and vertexes between the facets which can have catalytic properties that differ drastically from the ones on almost all other surfaces. For example, gold, which is rather inert as a bulk metal, has high reactivity in the form of nanosized clusters; the origin of their catalytic properties is still widely debated, however [14–16]. Research related to active sites is mainly limited to atomic steps on surfaces, and nanosized clusters have received much less attention [17–20]. If we want to understand the active sites on clusters, we have to know which cluster we are studying. Even then each cluster has several different active sites, and it is very difficult to know which of them is the most active. For this reason the computational approach, where precise sites can be studied, is very attractive.

Carbon monoxide (CO) is one of the most common carbon sources for the production of CNTs [9, 21–23]. This is due to the reasonable temperature range for CO disproportionation and the good CNT yield. Also the thermal self-decomposition of CO, leading to the precipitation of amorphous carbon and the formation of soot particles, can be avoided. The essential role of etching agents ( $\text{CO}_2$  and  $\text{H}_2\text{O}$ ) in CNT formation in the CVD process has been demonstrated [24]. These molecules are assumed to etch the amorphous carbon that can poison the catalyst particles needed for CO disproportionation, as well as for the CNT nucleation and growth. The important role of  $\text{H}_2$  in the CVD process has also been shown [31]. The  $\text{H}_2$  is believed to prevent the oxidation of the nanosized catalyst iron particles. Another possible role is in a reaction with  $\text{CO}_2$ :

$\text{CO}_2 + 2\text{H}_2 \rightarrow 2\text{H}_2\text{O} + \text{C}$ , which can provide additional carbon for the CNT synthesis. An earlier work has attempted to explain the effect of hydrogen in carbon fiber formation [27]. They suggested that the role of  $\text{H}_2$  is to limit or alleviate the carbon poisoning of the metal catalysts. More recently, it has been proposed that an intermediate, known as precursor soot, is involved in the formation of multi-walled carbon nanotubes (MWCNTs) and that  $\text{H}_2$  reduces the rate of dehydrogenation of the  $\text{CH}_x$  intermediates which form the soot [28].

To understand the chemistry of CNT formation on a catalytic particle several reactions need to be considered. In the ideal case, with only CO as precursor, these include the dissociation of CO,  $\text{CO}_2$  formation and release of C, the migration of carbon into the bulk of the catalyst, segregation of carbon to the surface, and the formation of graphene networks. These reactions are already quite complex and any additional new atom will make the reactions much more complex and thus a complete study with three species, like C, O, and H, is probably an unreachable goal with present computational resources. For these reasons, this study addresses the disproportionation of isolated CO molecules on iron nanoclusters ( $\text{CO}_{(\text{g})} + \text{CO}_{(\text{g})} \leftrightarrow \text{CO}_{2(\text{g})} + \text{C}_{(\text{s})}$ ) and the dissociation of  $\text{H}_2$ . The detailed chemistry on the surface of CNT catalyst transition metal nanoparticles is largely unknown, but it is believed that the adsorbed CO first dissociates,  $\text{CO}_{(\text{ads})} \rightarrow \text{C}_{(\text{ads})} + \text{O}_{(\text{ads})}$ , and then  $\text{O}_{(\text{ads})}$  reacts with an undissociated  $\text{CO}_{(\text{ads})}$  to form surface  $\text{CO}_{2(\text{ads})}$  (Fig. 1) [25, 26].

In order to understand the role of the cluster in this reaction, we have used first principles calculations to study the steps of this reaction on a small nanocluster. To our knowledge, there have been few ab initio studies of nanosized clusters [29, 30] and these works do not address any chemical



**Figure 1** Proposed CO disproportionation mechanism (Fe, grey; C, green; and O, red)

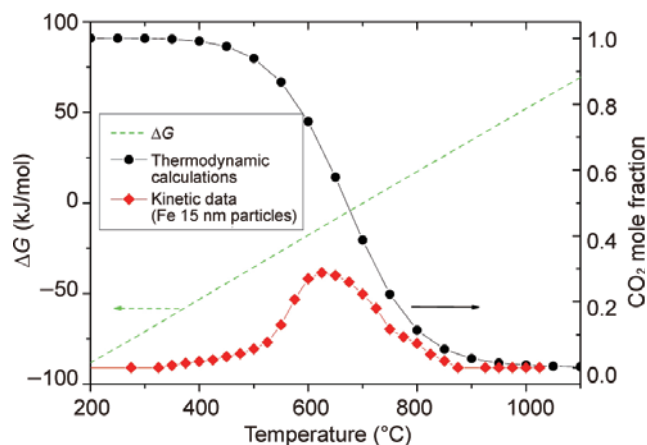


reactions. The nanoparticles have features that cannot be studied on stepped surfaces, and this work will complement existing studies of such active sites. The computational work is combined with experimental investigations of the same reaction on larger nanoclusters [31, 32].

## 1. Methods

The method for the production of single-walled carbon nanotubes (SWCNTs) has been described in detail elsewhere [31, 32]. In this method, pre-made catalyst particles of controlled size, concentration, and polydispersity are introduced under well defined synthesis conditions. Briefly, catalyst particles were produced from Fe vaporized from a resistively heated catalyst wire in a  $\text{H}_2/\text{N}_2$  (with a 7/93 molar ratio) flow ( $400 \text{ cm}^3/\text{min}$ ). Particles were formed and grown by vapor nucleation, condensation, and particle collision processes. Subsequently, the resulting particles were introduced into either a ceramic or stainless steel tubular reactor, mixed with a CO flow ( $400 \text{ cm}^3/\text{min}$ ) and heated to induce CNT formation. Reactor wall peak temperatures ranged from  $890^\circ\text{C}$  to  $1207^\circ\text{C}$ .

For electronic structure calculations, we have used density functional theory (DFT) with gradient correction and plane wave basis set within the super-cell approach (VASP code [33]). The adequacy of the functional form, pseudopotentials and plane wave basis set has been tested against the calculation of cluster geometry. Iron has a large magnetic moment and in the cluster geometry, the spin orientation is determined by subtle effects. Therefore, we have fully incorporated the effects of spin polarization and non-collinear magnetic moments in our calculations [34]. To our knowledge this is the first time that these effects have been included when substrate molecules or reactivity on nanoclusters is considered. Preliminary test calculations of the magnetic and structural properties of iron clusters up to  $\text{Fe}_{55}$  has been performed using  $E_{\text{Cut-off}}$  up to 270 eV with cubic super cell with size between  $17 \text{ \AA}$  and  $30 \text{ \AA}$ . This cluster is the best choice for our study because it is the largest one that we have been able to simulate whilst obtaining good results for the



**Figure 2** Temperature dependence of the Gibbs free energy for the CO disproportionation reaction. The equilibrium concentration of  $\text{CO}_2$  and kinetic data for CO disproportionation on the surface of 15 nm sized iron particles as a function of temperature are also plotted in the figure

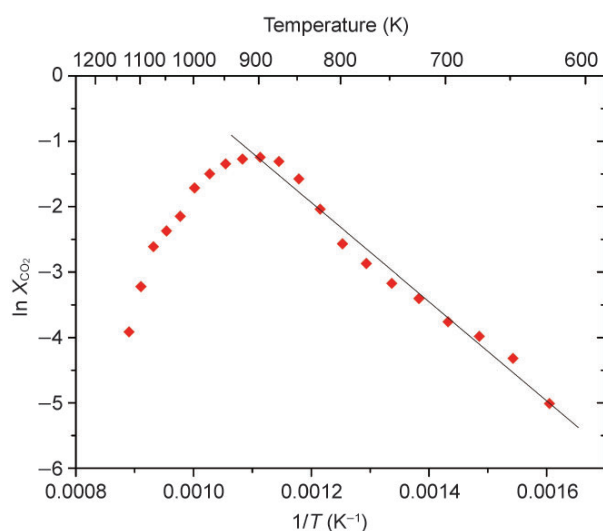
magnetic and geometric properties [34]. Good results were obtained with a super-cell size of  $21 \text{ \AA}$ , with  $E_{\text{Cut-off}} = 250 \text{ eV}$ . These settings are also used for the rest of calculations. The cluster has an icosahedral geometry and has a formation energy of  $-3.87 \text{ eV/atom}$ . The largest Fe–Fe distance was  $9.68 \text{ \AA}$ . The dipole moment was  $(-0.05; -0.03; 0.04) \text{ e l \AA}$ ; and the magnetic moment was  $(2.33; 0.48; 1.06) \text{ \mu B/atom}$ . The reactions barriers are more demanding to study and special transition state search algorithms, like climbing-image nudged elastic band (CI-NEB) [39], have been used for their determination. According to this algorithm, a set of  $3N$ -dimensional images of the adsorbate system are generated between the endpoint configurations. A harmonic interaction between adjacent images is added to ensure continuity of the path, thus mimicking an elastic band. An optimization of the band, involving the minimization of the forces acting on the images, brings the band to the minimum energy path.

## 2. Results

A gas-phase process for formation of SWCNTs, based on thermal decomposition of iron pentacarbonyl or ferrocene in the presence of carbon monoxide (CO), was investigated in ambient pressure laminar flow reactors in the temperature range  $600\text{--}1300^\circ\text{C}$  [31]. *In situ* sampling carried out at  $1000^\circ\text{C}$  showed that the growth of SWCNTs occurred from individual

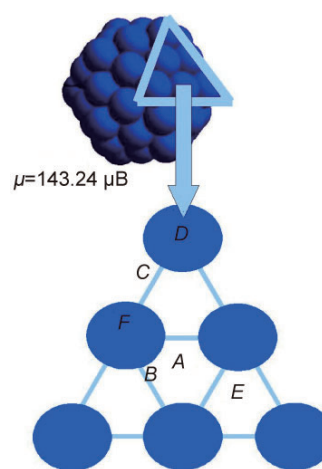
metal particles in the heating section of the furnace in the temperature range 891–928 °C, during which the growth rate was estimated to exceed 2  $\mu\text{m/s}$  [32]. Kinetic investigations of the CO disproportionation reaction were performed in a horizontal quartz tube at a heating rate of 5 °C/min from room temperature. A silica substrate with deposited 15 nm sized iron particles was placed inside the tube. The CO flow rate through the tube was maintained at 8.3  $\text{cm}^3/\text{min}$ . The concentration of the gaseous product ( $\text{CO}_2$ ) was monitored on-line by infrared (IR). Investigations using 15 nm sized iron particles show appreciable reaction rates in the temperature interval from 470 to 820 °C with a maximum rate at about 625 °C (Fig. 2). The region of  $\text{CO}_2$  concentration increase from about 325 °C to about 600 °C is the kinetic region, where the rate of the CO disproportionation reaction can be measured. At temperatures above 625 °C, this reaction is limited thermodynamically as can be seen from the values of Gibbs free energy,  $\Delta G$ . Since the disproportionation reaction rate is proportional to the concentration of  $\text{CO}_2$ , the activation energy of the reaction can be found using an Arrhenius dependence:  $X_{\text{CO}_2} = k_0 \exp(-E_a/RT)$ , where  $X_{\text{CO}_2}$  is the carbon dioxide mole fraction. Plotting the kinetic region in the coordinates of  $\ln X_{\text{CO}_2}$  versus  $1/T$ , one gets  $k_0 = 7.16091$  and  $E_a = 63031 \text{ J/mol} = 0.65 \text{ eV}$  as an energy barrier for the CO disproportionation reaction (Fig. 3).

Theoretical calculations (VASP code [33]) have



**Figure 3** Kinetic region of the CO disproportionation reaction fitted with a straight line (Arrhenius plot)

been used to study the geometry of a 1 nm size  $\text{Fe}_{55}$  cluster and to identify the stationary points of the species involved in the reactions on it. In order to understand the reaction mechanism, several reaction paths have been studied. Our results suggest that the energetically most favorable cluster has an icosahedral symmetry. In this conformation, the 55-atom cluster has three atomic layers, one atom in the center, a 12-atom inner icosahedron, and a 42-atom outer icosahedron. Our results are in good agreement with the density functional based study of Kohler et al. [34]. The geometries and bonding energies of the adsorbates ( $\text{CO}$ ,  $\text{CO}_2$ ,  $\text{H}_2$ , and atomic C/O) were calculated using the same computation approach used for the icosahedral  $\text{Fe}_{55}$ . The geometry optimization was started from the high symmetry adsorption sites on one of the cluster faces. The high symmetry adsorption sites are the top site ( $T_s$ ), the octahedral hollow site ( $H_s$ ), and the bridge sites ( $B_s$ ). The  $T_s$  and  $B_s$  sites exist both on the cluster faces and on its edges/Vertexes.  $H_s$  sites are on the triangular faces of the icosahedral  $\text{Fe}_{55}$ : A and E in plane hexagonal close packed (hcp) site (Fig 4).  $B_s$  sites are located in on the facet (plane bridge site) like B (Fig 4); or on the edge bridge site (C in Fig 4);  $T_s$  site are vertex top sites, D and F, on edge top site (Fig. 4). The adsorption energy ( $E_{\text{ads}}$ ) is defined as  $E_{\text{ads}} = E_{\text{Fe}_{55} + \text{mol}} - (E_{\text{Fe}_{55}} + E_{\text{mol}})$ , where  $E_{\text{Fe}_{55} + \text{mol}}$  is the total energy of the cluster with adsorbates,  $E_{\text{Fe}_{55}}$  is the total energy of the



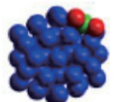
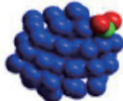
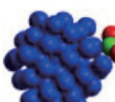
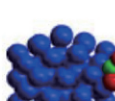
**Figure 4** High symmetry adsorption sites on a triangular face of the icosahedral  $\text{Fe}_{55}$ : A, in plane hexagonal close packed (hcp) site; B, in plane bridge site; C, on edge bridge site; D, on vertex top site; E, in plane hcp site; and F, on edge top site



$\text{Fe}_{55}$  cluster, and  $E_{\text{mol}}$  is the total energy of free  $\frac{1}{2} \text{O}_2$ ,  $\frac{1}{2} \text{H}_2$ , CO,  $\text{CO}_2$  molecules or C atom. Thus, the more negative is  $E_{\text{ads}}$ , the stronger is the adsorption.

For CO, geometries with both the molecular bond coplanar (=) or perpendicular (p) to the cluster surface were considered. The total magnetic moment of each structure is reported in the figures and tables. Only the perpendicular adsorption with the molecule adsorbed through the carbon is stable. During the optimization, CO spontaneously moved to the edge of the face (*D* and *F* sites in Fig. 4). The most favored adsorption site for CO is the  $T_s$  site on the vertex of the cluster (*F* in Fig. 4) having  $E_{\text{ads}} = -1.83$  eV, followed by the *D* site with 0.40 eV higher energy. On the cluster *F* and *D* sites, the C–O bond is elongated to 1.181 Å and 1.179 Å, respectively (cf. 1.144 Å for free CO). These results can be compared with the reported binding energies and stable geometries for CO on different Fe surfaces at low coverage ( $\theta = 0.25$ ). For the stepped Fe (211) surface, DFT calculations by Borthwick et al. [46] propose the p-CO interaction on the step of the surface as the most stable geometry ( $E_{\text{ads}} = -1.73$  eV, and C–O bond length 1.17 Å). This is in good agreement with the lowest values,  $-1.81$  eV and  $-1.43$  eV, reported here. DFT calculations of Chen et al. [35] gave  $E_{\text{ads}} = -1.65$  eV for the p-CO adsorption on the  $T_s$  site of the Fe (111) surface. A similar on-top adsorption geometry has also been obtained on the less open Fe (110) and Fe (100) surfaces [36, 45] with respective adsorption energies of  $-2.00$  eV and  $-1.392$  eV. These results suggest that at low coverage, the CO adsorption on the  $\text{Fe}_{55}$  cluster is comparable to previously studied Fe surfaces.

A systematic study of the stability of  $\text{CO}_2$  on the vertex and edge of the cluster was considered as well as on the planar faces, through a series of geometry optimizations starting from several different molecular orientations. Among these possible initial geometries, only four stable conformations were found and the  $\text{CO}_2$  geometry was distorted. The adsorption energies ( $-0.39 \text{ eV} < E_{\text{ads}} < -0.11 \text{ eV}$ ) and corresponding geometries are reported in Fig. 5. The minimum energy configuration ( $E_{\text{ads}} = -0.39$  eV) corresponds to the bridge site with one of the molecular bonds lying on the cluster surface (labeled as  $\text{CO}_2\_B1$  in Fig. 5). In this configuration, the  $\text{CO}_2$  molecule is bent

Adsorbate	Adsorption site	$E_{\text{ads}}$ (eV)	$\mu$ ( $\mu\text{B}$ )	Geometry	
$\text{CO}_2$ (molecule)	$B$	$\text{CO}_2\_B1$	−0.39	143.68	
		$\text{CO}_2\_B2$	−0.29	143.61	
	$F$	$\text{CO}_2\_F1$	−0.24	143.39	
		$\text{CO}_2\_F2$	−0.11	143.53	

**Figure 5** Stable geometries of  $\text{CO}_2$  on  $\text{Fe}_{55}$  (Fe, blue; C, green; and O, red)

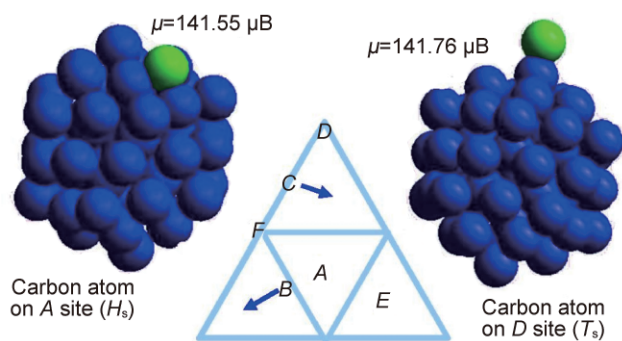
(molecular angle =  $152^\circ$ ), with the carbon–oxygen bond closer to the surface at a distance of 1.90 Å and elongated to 1.36 Å (cf. 1.16 Å for the free molecule). The other C–O bond has a length of 1.31 Å. Adsorption is also possible on *F* sites. In this case, the molecule has  $C_{2v}$  symmetry and the molecular plane is upright to the cluster and parallel or perpendicular to the edge direction. These latter cases are local minimum configurations corresponding to weaker bonding, with  $E_{\text{ads}} = -0.24$  eV. In this configuration (labeled as  $\text{CO}_2\_F1$  in Fig. 5), the  $\text{CO}_2$  molecule is bent, with the carbon atom closer to the surface at about 1.92 Å and the C–O bond lengths are equal to 1.28 Å. The adsorption of  $\text{CO}_2$  on various transition metals has been the subject of many investigations. The geometries we obtained are in accordance with the results suggested in the previous literature [37].

In order to study CO dissociation and  $\text{CO}_2$  formation with the (CI-NEB) method [38, 39], we need to investigate the preferred C and O coadsorption on the cluster. First, all possible adsorption sites were considered, and only three stable sites for individually adsorbed C and O were

located. Both of the atoms show the same behavior (Fig. 6):  $H_s$  or  $T_s$  are the only stable sites and from the bridge sites, the adsorbate migrates to the nearest hollow site. Table 1 displays the adsorption energies of isolated C or O atoms on the high-symmetry sites; for comparisons remember that the reference energy for carbon is the isolated atom, and that for the oxygen is half of the  $O_2$  molecule. We found that the  $H_s$  sites ( $E$  in Fig. 6) are the most favorable ones with  $E_{\text{ads}}(\text{O}) = -4.27$  eV and  $E_{\text{ads}}(\text{C}) = -7.96$  eV. The observed  $\text{Fe}_3\text{C}$  structures, corresponding to carbon adsorption on the  $H_s$  sites, were also previously observed on different Fe surfaces [40–42].

**Table 1** Binding energies for the main adsorption geometries of  $\text{H}_2$ , atomic oxygen and atomic carbon on  $\text{Fe}_{55}$

Adsorption site	Binding energy (eV)		
	$\text{H}_2$	O	C
A	+0.09	-3.87	-7.91
B	-0.36	-4.23	-7.99
C	+0.28	-4.26	-7.98
D	-0.03	-2.96	-5.83
E	+0.26	-4.27	-7.96
F	-0.03	-4.23	-6.44



**Figure 6** Schematic representation of the stable geometries of atomic carbon on  $\text{Fe}_{55}$ . Atomic oxygen shows similar behavior. (Fe, blue; and C, green)

Next, we placed both C and O atoms near each other on the same face of the cluster. All the possible coadsorption combinations were considered, but only six geometries (reported in Fig. 7) have exothermic dissociative chemisorption ( $\Delta E$ ), defined as  $\Delta E = (E_{\text{Fe}_{55}/\text{C}+\text{O}} - (E_{\text{Fe}_{55}} + E_{\text{CO}})) / N_{\text{CO}}$ , where  $E_{\text{Fe}_{55}/\text{C}+\text{O}}$  is the total energy of the cluster with atomic C/O co-adsorbates and  $N_{\text{CO}}$  is the number of free CO molecules. Interestingly, it seems that having an O

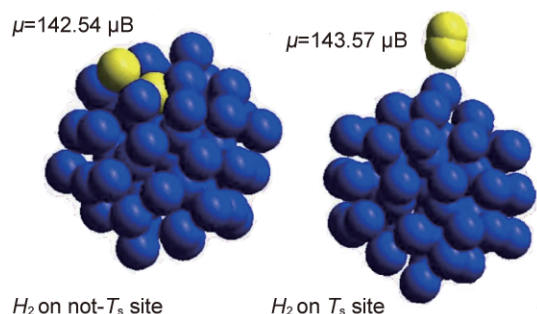
atom nearby pushes the C atom onto another face of the cluster. We also found that for the most stable geometries (indicated as C\_O\_2 and C\_O\_3 in Fig. 7), the O atom remains on the surface while the C atom is embedded into cluster, indicating C atoms may indeed diffuse into the bulk. This carbide-like behavior was suggested in the previous literature [43]. This behavior is also in line with previous Raman ellipsometry spectroscopy observation [49] of surfaces. In these studies, Watanabe et al. found that CO dissociation produces a thin  $\text{Fe}_3\text{O}_4$  layer on Fe surfaces, which blocks a large number of C atoms in the metallic phase.

Adsorbate	Adsorption site	$\Delta E$ (eV)	$\mu$ ( $\mu\text{B}$ )	Geometry
C (atom) + O (atom) on the same face	C_O_1 O on A site C on E site	-0.32	141.01	
	C_O_2 O on E site C on C site	-2.44	141.06	
C (atom) + O (atom) on different face	C_O_3 O on A site C on A site	-2.50	141.03	
	C_O_4 O on E site C on E site	-1.97	141.77	
	C_O_5 O on E site C on E site	-1.78	141.40	
	C_O_6 O on E site C on E site	-0.99	141.90	

**Figure 7** Stable geometries of C+O on  $\text{Fe}_{55}$  showing exothermic dissociative chemisorption (Fe, blue; C, green; and O, red)

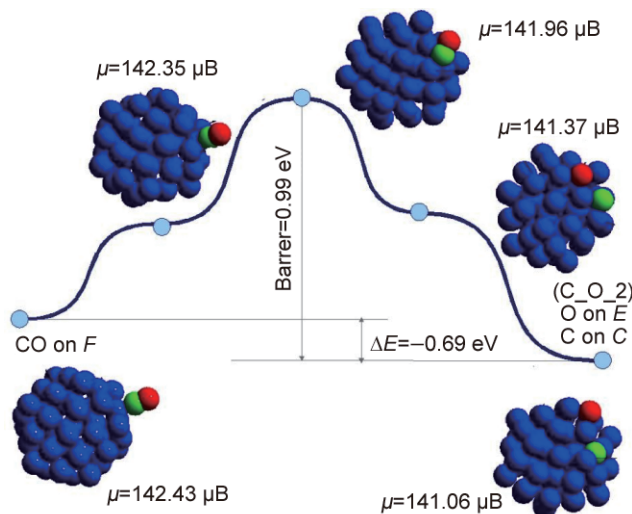
The stability of  $\text{H}_2$  on the cluster was also considered. The resulting binding energies for  $\text{H}_2$  adsorption perpendicular to the cluster surface are reported in Table 1. During the geometry optimization starting from the top sites, the molecule keeps its position in a metastable state, but from the other sites  $\text{H}_2$  spontaneously dissociates and one

adsorbate migrates below the surface. Schematic representation of the main adsorption geometries is reported in Fig. 8. In this work we will not study further the reactions of hydrogen and other species, but the atomic hydrogen will react very easily with O, C, and larger C clusters.

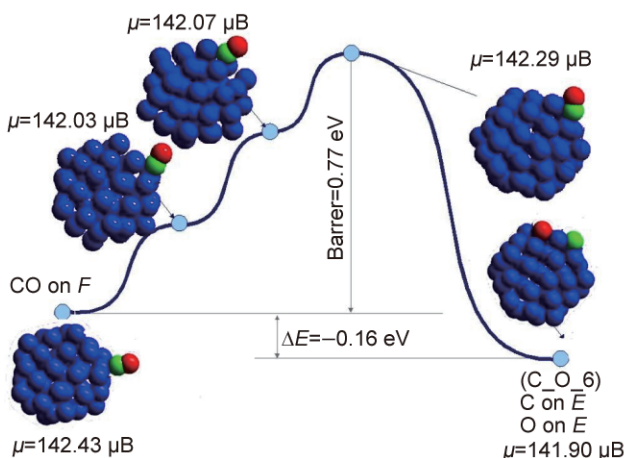


**Figure 8** The main adsorption geometries of  $H_2$  on  $Fe_{55}$  (Fe, blue; and H, yellow)

Based on CO adsorption and C/O co-adsorption data, CO dissociation paths from the most stable CO adsorption geometries (on  $F$  sites and on  $D$  sites) to the best C/O coadsorption structures (Fig. 7) were mapped out with the CI-NEB method. Several minimum energy paths (MEPs) were calculated and the barriers obtained vary between 0.77 eV and 3.27 eV. Here we focus on the two lowest barrier paths,  $CO_G \rightarrow C/O_{C_{O_2}}$  (barrier 0.99 eV, shown in Fig. 9) and  $CO_G \rightarrow C/O_{C_{O_4}}$  (barrier 0.77 eV, shown in Fig. 10). The path  $CO_G \rightarrow C/O_{C_{O_2}}$  is thermodynamically the most favored one of the studied paths. This reaction is exothermic by 0.69 eV. In the transition state, the C–O bond is elongated to 1.25 Å, with the O atom close to a bridge site (Fig. 9). This type of transition state is quite common for CO dissociation on transition metal surfaces [44, 45, 48]. In the final structure, the C atom becomes embedded into the cluster surface, whereas the O atom is on the surface. Although the path  $CO_G \rightarrow C/O_{C_{O_2}}$  has the largest  $\Delta E$ , the  $CO_G \rightarrow C/O_{C_{O_2}}$  path is kinetically the most favored one. This reaction is exothermic by 0.16 eV. The CO first moves to the edge of the face, reorients itself, and then goes through a transition state where CO tilts on the vertex of the cluster. Early temperature-programmed desorption (TPD) and time-resolved electron energy loss spectroscopy (TREELS) studies conducted by Whitman et al. [50] determined



**Figure 9** Calculated minimum energy paths for CO dissociation on a facet of  $Fe_{55}$ . Several intermediate structures along the path are shown (Fe, blue; C, green; and O, red)



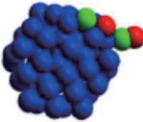
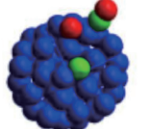
**Figure 10** Calculated minimum energy paths for CO dissociation on a step of  $Fe_{55}$ . Several intermediate structures along the path are shown (Fe, blue; C, green; and O, red)

the kinetic data for CO adsorption and reaction on  $Fe(111)$ . It was observed that at a monolayer coverage of 0.312, in addition to dissociation, some of the adsorbed CO desorbs with a barrier of ~1.4 eV. This can be taken as the upper limit for the CO dissociation barrier on this surface. If the predicted minimum MEPs for CO dissociation are compared with those from calculations for other surfaces, it can be seen that the paths are similar and but the presence of the vertex lowers the barrier of the dissociation reaction. The activation barrier of 0.77 eV obtained during our calculation is significantly lower than most of the reported barriers for different

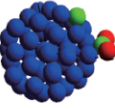
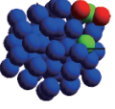


surfaces (even the lowest case, 0.78 eV [46], is very close to that in our calculations), suggesting that the  $\text{Fe}_{55}$  is much more reactive towards CO dissociation than most of the stepped Fe surfaces. It seems that the ability of the metal atoms to activate reactants changes substantially as their coordination number is reduced, both at the atomic steps and the vertexes. The results show that the catalytically most active sites are the vertexes where the molecules also have a relatively low coordination number.

We consider a mechanism for the CO oxidation during the CVD synthesis that involves the reaction of the molecule with an oxygen atom adsorbed on the surface. In order to study the formation of  $\text{CO}_2$  from CO and its fragment adsorbed on the cluster, the CO/O/C and  $\text{CO}_2$ /C co-adsorption geometries have been studied. The most stable geometries are indicated in Figs. 11 and 12, respectively. Taking  $\text{CO}_2\text{C}_1$  (Fig. 11) as the starting point, possible paths to  $\text{CO}_2\text{C}_1$  (Fig. 12) can be easily conceived. The corresponding structures of all transition states and intermediates are depicted in Fig. 13 together with the energy profile. The formation energy obtained is  $-0.37$  eV, with a barrier of 1.13 eV. This path presents an elementary rearrangement of atoms and molecules on the surface. A more complex reaction path was found taking  $\text{CO}_2\text{C}_2$  (Fig. 11) as the starting point and  $\text{CO}_2\text{C}_2$  (Fig. 12) as the final point of the path. This path has a formation energy of  $+0.88$  eV and barrier of 1.08 eV. Considering that the particle morphology and the reaction conditions are different, the calculated data are in reasonable agreement with

Adsorbate	Adsorption site				$\Delta E$ (eV/CO molecule)	$\mu$ ( $\mu\text{B}$ )	Geometry
	Label	CO molecule	C atom	O atom			
CO (molecule) + O (atom) + C (atom)	$\text{CO}_2\text{C}_1$	F	E	A	$-0.78$	139.10	
	$\text{CO}_2\text{C}_2$	F	E	E	$-1.80$	139.44	

**Figure 11** Stable geometries of  $\text{CO} + \text{C} + \text{O}$  on  $\text{Fe}_{55}$  (Fe, blue; C, green; and O, red)

Adsorbate	Adsorption site			$\Delta E$ (eV/ $\text{CO}_2$ molecule)	$\mu$ ( $\mu\text{B}$ )	Geometry
	Label	$\text{CO}_2$ molecule	C atom			
$\text{CO}_2$ (molecule) + C (atom)	$\text{CO}_2\text{C}_1$	B	E	$-0.97$	140.55	
	$\text{CO}_2\text{C}_2$	F	E	$-1.36$	140.24	

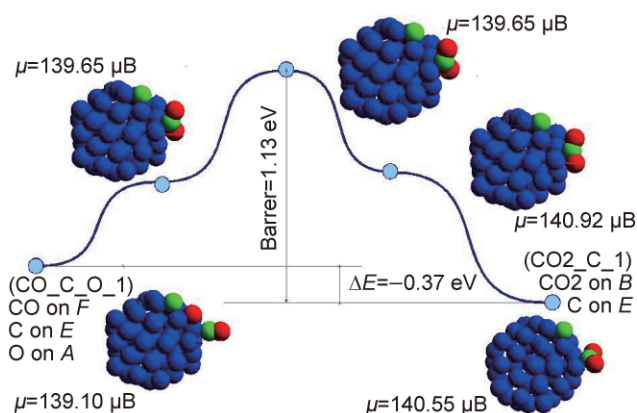
**Figure 12** Stable geometries of  $\text{C} + \text{CO}_2$  on  $\text{Fe}_{55}$  (Fe, blue; C, green; and O, red)

the experimental results of Li et al. [26] The calculated barriers are, furthermore, in good agreement with the experimental results for the  $\text{Fe}_2\text{O}_3$  catalysts used in modern industrial applications ( $\sim 0.90$  eV) [47].

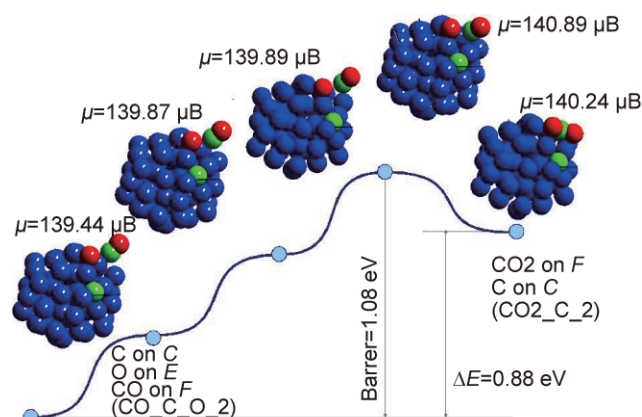
We can also compare the calculated results with our experimental results. The DFT results above show that CO disproportionation on a nanosized iron cluster is a multistep process in which the rate limiting step is the formation of  $\text{CO}_2$  from O and CO. We studied different reaction channels, and the lowest energy channel showed a barrier of 1.08 eV that is similar to the value obtained from the kinetic investigations, with a barrier of 0.65 eV for the overall process. The difference between the calculated and experimental results is reasonable and it can be justified considering that the cluster geometry and dynamics differ markedly in the two cases. The computed nanoparticle is very small and regular and thus it has few active sites. In this sense it is probably rather unreactive as a nanocluster. The larger clusters relevant to the experiments will have many more active sites and some of these can have lower activation barriers. Also, our calculations are at 0 K, so Fe atom diffusion is not taken into account, whereas the experiments are carried out around 900 K where the Fe atoms are very mobile. The iron mobility can also generate new active sites that are very difficult to model with static calculations. Even although the difference between the experimental and simulated clusters is rather large, we want to emphasize that the calculations have provided new information about the reactivity of nanoclusters and comparisons to real-world systems are more relevant at lower temperatures.







**Figure 13** Calculated minimum energy paths for  $\text{CO}_2$  formation on  $\text{Fe}_{55}$  starting from CO and co-adsorbed atomic C/O. Several intermediate structures along the path are shown (Fe, blue; C, green; and O, red)



**Figure 14** Calculated minimum energy paths for  $\text{CO}_2$  formation on  $\text{Fe}_{55}$  starting from CO and co-adsorbed atomic C/O. Several intermediate structures along the path are shown (Fe, blue; C, green; and O, red)

### 3. Conclusions

We have presented an *ab initio* study of CO adsorption and dissociation,  $\text{H}_2$  dissociation, and  $\text{CO}_2$  formation on a small iron nanoparticle. To our knowledge, this is the first computational study of a chemical reaction on a real nanosized iron particle. Even more importantly, we have studied active sites that have not been studied before, the facets and vertexes. Overall, we see the small clusters as new and important model systems for nanocatalyst studies. They exhibit new active sites that complement those on stepped surfaces studied in other calculations. When the computational capacity permits, larger clusters can also be studied. At present, detailed chemical reactions can be

studied for clusters with a size of 50–100 atoms. Another novel aspect of our study is the use of spin polarization with non-collinear magnetic moments.

In our calculations, the  $\text{Fe}_{55}$  cluster is found to have high catalytic activity. We have observed one important effect: the increased reactivity of low coordinated molecules on the vertexes. Such sites are particularly abundant on small nanosized particles, and they very likely explain their catalytic activity. We have been able to show that the reaction barriers for both CO dissociation and  $\text{CO}_2$  formation are relatively low and an exothermic path for surface bound elemental carbon has been found. Carbon atom formation is the initial step in CNT formation. Under CVD conditions, 1.5–2.5 nm sized clusters are the CNT catalysts. These clusters probably have a similar geometry to what we have studied here, and we believe that in this case the vertexes are also the catalytically active sites.

In this work, we have studied the carbon disproportionation reactions of CO and  $\text{H}_2$  dissociation. In further studies, we will focus on subsurface carbon, carbon–carbon, hydrogen–oxygen and hydrogen–carbon reactions on the nanoparticle. These studies will give information about carbide formation and the possible role of hydrogen. We see that it is very challenging to study the full process of CNT formation with *ab initio* methods. Thus, less detailed methods, like tight-binding, could also be used, but for some of the early reaction steps the *ab initio* calculations can give very reliable information. Also such calculations will be an excellent reference for the less detailed calculations.

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